

## Completeness of Bethe's states for the generalized **XXZ** model

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1997 J. Phys. A: Math. Gen. 30 1209

(<http://iopscience.iop.org/0305-4470/30/4/022>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.112

The article was downloaded on 02/06/2010 at 06:12

Please note that [terms and conditions apply](#).

# Completeness of Bethe's states for the generalized $XXZ$ model

Anatol N Kirillov<sup>†</sup> and Nadejda A Liskova<sup>‡</sup>

<sup>†</sup> Steklov Mathematical Institute, Fontanka 27, St Petersburg, 191011, Russia

<sup>‡</sup> St Petersburg Institute of Aviation Instruments, Gertzena 67, St Petersburg, 190000, Russia

Received 21 May 1996, in final form 28 August 1996

*Dedicated to the memory of Ansgar Schnizer*

**Abstract.** We study the Bethe ansatz equations for a generalized  $XXZ$  model on a one-dimensional lattice. Assuming the string conjecture we propose an integer version for vacancy numbers and prove a combinatorial completeness of Bethe's states for a generalized  $XXZ$  model. We find an exact form for the inverse matrix related with vacancy numbers and compute its determinant. This inverse matrix has a tridiagonal form, generalizing the Cartan matrix of type  $A$ .

## 1. Introduction

An integrable generalization of spin- $\frac{1}{2}$  Heisenberg  $XXZ$  model to arbitrary spins was given, for example, in [KR2]. As a matter of fact, a spectrum of the generalized  $XXZ$  model is described by the solutions  $\{\lambda_i\}$  to the following system of equations ( $1 \leq j \leq l$ ):

$$\prod_{a=1}^N \frac{\sinh \frac{1}{2}\theta(\lambda_j + 2is_a)}{\sinh \frac{1}{2}\theta(\lambda_j - 2is_a)} = \prod_{\substack{k=1 \\ k \neq j}}^l \frac{\sinh \frac{1}{2}\theta(\lambda_j - \lambda_k + 2i)}{\sinh \frac{1}{2}\theta(\lambda_j - \lambda_k - 2i)}. \quad (1.1)$$

Here  $\theta$  is an anisotropy parameter,  $s_a$ ,  $1 \leq a \leq N$ , are the spins of atoms in the magnetic chain and  $l$  is the number of magnons over the ferromagnetic vacuum.

The main goal of our paper is to present a computation the number of solutions to system (1.1) based on the so-called string conjecture (see, e.g., [TS], [KR1]). In spite of the well known fact that solutions of (1.1) do not have, in general, a 'string nature' (see, e.g., [EKK]), we prove that the string conjecture gives a correct answer for the number of solutions to the system of equations (1.1). Note that a combinatorial completeness of Bethe's states for the generalized  $XXX$  Heisenberg model was proved in [K1] and appears to be a starting point for numerous applications to combinatorics of Young tableaux and representation theory of symmetric and general linear groups (see, e.g., [K2]).

## 2. Analysis of the Bethe equations

Let us consider the  $XXZ$  model of spins  $s_1, \dots, s_k$  interacting on a one-dimensional lattice with the each spin  $s_i$  repeated  $N_i$  times. In the standard  $XXZ$  model all spins  $s_i$  are equal to  $\frac{1}{2}$ . Let  $\Delta$  be the anisotropy parameter (see, e.g., [TS], [KR2]). We assume that  $0 < \Delta < 1$ .

Let us pick out a real number  $\theta$  such that  $\cos \theta = \Delta$ ,  $0 < \theta < \frac{\pi}{2}$ , and denote

$$p_0 = \frac{\pi}{\theta} > 2.$$

Each spin  $s$  has a ‘parity’  $v_{2s}$  which is equal to plus or minus one.

Bethe vectors  $\psi(x_1, \dots, x_l)$  for  $XXZ$  model are parametrized by  $l$  complex numbers  $x_j \pmod{2p_0i}$  ( $l \leq 2s_1N_1 + \dots + 2s_kN_k$ ), which satisfy the following system of transcendental equations (Bethe’s equations),

$$\prod_{m=1}^k (-1)^{N_m v_{2s_m}} \left( \frac{\sinh \frac{1}{2}\theta(x_\alpha + \eta_m - i(2s_m + \frac{1}{2}(1 - v_{2s_m})p_0))}{\sinh \frac{1}{2}\theta(x_\alpha + \eta_m + i(2s_m + \frac{1}{2}(1 - v_{2s_m})p_0))} \right)^{N_m} = - \prod_{j=1}^l \frac{\sinh \frac{1}{2}\theta(x_\alpha - x_j - 2i)}{\sinh \frac{1}{2}\theta(x_\alpha - x_j + 2i)} \tag{2.1}$$

where  $\alpha = 1, \dots, l$  and  $\{\eta_m\}$  are some fixed real numbers, and the non-degeneracy conditions, the norm of the Bethe’s vectors  $\psi$  are not equal to zero, apply.

Solutions to the system (2.1) are considered modulo  $2p_0i\mathbb{Z}$ , because  $\sinh(\frac{1}{2}\theta x)$  is a periodic function with the period  $2p_0i$ . Asymptotically for  $N_m \rightarrow \infty$ ,  $1 \leq m \leq k$  and finite  $l$  the solutions to the system (2.1) create the strings. The strings are characterized by the common real abscissa, which is called the string centre, the length  $n$  and parity  $v_n$ . Centres of even strings are located on the line  $\text{Im } x = 0$  (and  $v_n = +1$ ), those of odd strings are located on the line  $\text{Im } x = p_0$  (and  $v_n = -1$ ). A string of length  $n$  and parity  $v_n$  consists of  $n$  complex numbers  $x_{\beta,j}^n$  of the following form,

$$x_{\beta,j}^n = x_\beta^n + i \left( n + 1 - 2j + \frac{1 - v_n}{2} p_0 + O(\exp(-\delta N)) \right) \pmod{2p_0i} \tag{2.2}$$

where  $\delta > 0$ ,  $j = 1, \dots, n$ ,  $x_\beta^n \in \mathbb{R}$ .

A distribution of numbers  $\{x_j\}$  on strings is called a configuration. Each configuration can be parametrized by the filling numbers  $\{\lambda_n\}$ , where  $\lambda_n$  is equal to the number of strings with length  $n$  and parity  $v_n$ . Each real solution of the system (2.1) (modulo  $2p_0i$ ), corresponds to an even string of length 1. Configuration parameters  $\{\lambda_n\}$ ,  $n \geq 1$ , satisfy the following conditions:  $\lambda_n \geq 0$ ,  $\sum_{n \geq 1} n\lambda_n = l$ . The system (2.1) can be transformed into that for real numbers  $x_\beta^n$  for each fixed configuration. To get such a system, let us calculate the scattering phase  $\theta_{n,m}(x)$  of the string length  $n$  on that of length  $m$ . By definition

$$\exp(-2\pi i \theta_{n,m}(x)) = \prod_{j=1}^n \prod_{k=1}^m \frac{\sinh \frac{1}{2}\theta(x_{\alpha,j}^n - x_{\beta,k}^m - 2i)}{\sinh \frac{1}{2}\theta(x_{\alpha,j}^n - x_{\beta,k}^m + 2i)} \quad x := x_\alpha^n - x_\beta^m. \tag{2.3}$$

From the formulae

$$\text{Im} \log \left( \frac{\sinh(\lambda + ai)}{\sinh(\mu + bi)} \right) = \arctan(\tanh \mu \cdot \cot b) - \arctan(\tanh \lambda \cdot \cot a) \tag{2.4}$$

where  $a, b, \lambda, \mu \in \mathbb{R}$ , it follows

$$\begin{aligned} -\pi \theta_{n,m}(x) &= \sum_{j=1}^n \sum_{k=1}^m \arctan(\tanh \frac{1}{2}\theta x \cot \frac{1}{2}\theta(n - m - 2j + 2k + \frac{1}{2}(v_m - v_n)p_0)) \\ &= \arctan(\tanh \frac{1}{2}\theta x \cdot \cot \frac{1}{2}\theta(m + n + \frac{1}{2}(1 - v_n v_m)p_0)) \\ &\quad + \arctan(\tanh \frac{1}{2}\theta x \cdot \cot \frac{1}{2}\theta(|n - m| + \frac{1}{2}(1 - v_n v_m)p_0)) \\ &\quad + 2 \sum_{s=1}^{\min(n,m)-1} \arctan(\tanh \frac{1}{2}\theta x \cdot \cot \frac{1}{2}\theta(|n - m| + 2s + \frac{1}{2}(1 - v_n v_m)p_0)). \end{aligned}$$

Now let us consider the limit of  $\theta_{n,m}(x)$  when  $x \rightarrow \infty$ . Note that for  $x \rightarrow \infty$ , we have  $\tanh \frac{1}{2}\theta x \rightarrow 1$ , and  $\arctan(\cot z) = -\pi((z/\pi))$ , if  $z/\pi \notin \mathbb{Z}$ , where  $((z))$  is the Dedekind function:

$$((z)) = \begin{cases} 0 & \text{if } z \in \mathbb{Z} \\ \{z\} - \frac{1}{2} & \text{if } z \notin \mathbb{Z} \end{cases}$$

and  $\{z\} = z - [z]$  is the fractional part of  $z$ . Then

$$\begin{aligned} \theta_{n,m}(\infty) &= \left( \left( \frac{n+m}{2p_0} + \frac{1-v_n v_m}{4} \right) \right) + \left( \left( \frac{|n-m|}{2p_0} + \frac{1-v_n v_m}{4} \right) \right) \\ &\quad + 2 \sum_{l=1}^{\min(n,m)-1} \left( \left( \frac{|n-m|+2l}{2p_0} + \frac{1-v_n v_m}{4} \right) \right). \end{aligned}$$

Let us define

$$\begin{aligned} \Phi_{n,m}(\lambda) &= -\frac{1}{2\pi i} \sum_{j=1}^n \log \frac{\sinh \frac{1}{2}\theta(x_{\alpha,j}^n + \eta - mi - \frac{1}{2}(1-v_m)p_0 i)}{\sinh \frac{1}{2}\theta(x_{\alpha,j}^n + \eta + mi + \frac{1}{2}(1-v_m)p_0 i)} \\ \lambda &:= x_{\alpha}^n + \eta \quad \eta \in \mathbb{R} \end{aligned}$$

then

$$\begin{aligned} \Phi_{n,m}(\lambda) &= -\frac{1}{2\pi} \sum_{j=1}^n 2 \cdot \arctan(\tanh \frac{1}{2}\theta x \cot \frac{1}{2}\theta(n-m+1-2j + \frac{1}{2}(v_m-v_n)p_0)) \\ &= \frac{1}{\pi} \sum_{l=1}^{\min(n,m)} \arctan(\tanh \frac{1}{2}\theta x \cot \frac{1}{2}\theta(|n-m|+2l-1 + \frac{1}{2}(1-v_n v_m)p_0)) \end{aligned}$$

and, consequently,

$$\Phi_{n,m}(\infty) = \sum_{l=1}^{\min(n,m)} \left( \left( \frac{|n-m|+2l-1}{2p_0} + \frac{1-v_n v_m}{4} \right) \right). \tag{2.5}$$

Now let us continue the investigation of system (2.1). Multiplying the equations of the system (2.1) along the string  $x_{\alpha,j}^n$  and taking the logarithm result, one can obtain the following system on real numbers  $x_{\alpha}^n, \alpha = 1, \dots, \lambda_n$ :

$$\sum_m \Phi_{n,2s_m}(x_{\alpha}^n + \eta_m) N_m = Q_{\alpha}^n + \sum_{(\beta,m) \neq (\alpha,n)} \theta_{n,m}(x_{\alpha}^n - x_{\beta}^m) \quad \alpha = 1, \dots, \lambda_n. \tag{2.6}$$

Integer or half-integer numbers  $Q_{\alpha}^n, 1 \leq \alpha \leq \lambda_n$ , are called quantum numbers. They parametrize—according to the string conjecture [TS], [FT], [K1]—the solutions to the system (2.1). Admissible values of quantum numbers  $Q_{\alpha}^n$  are located in the symmetric interval  $[-Q_{\max}^n, Q_{\max}^n]$  and appear to be an integer or half-integer in accordance with that of  $Q_{\max}^n$ .

### 3. Calculation of vacancy numbers

Following [TS], we will assume that there are two types of length 1 string, namely even and odd types. If the length  $n$  of a string is greater than 1, then  $n$  and parity  $v_n$  satisfy the following conditions:

$$v_n \cdot \sin(n-1)\theta > 0 \tag{3.1}$$

$$v_n \cdot \sin(j\theta) \sin(n-j)\theta > 0 \quad j = 1, \dots, n-1. \tag{3.2}$$

Condition (3.1) may be rewritten equivalently as

$$v_n = \exp\left(\pi i \left[\frac{n-1}{p_0}\right]\right)$$

and (3.2) as

$$\left[\frac{j}{p_0}\right] + \left[\frac{n-j}{p_0}\right] = \left[\frac{n-1}{p_0}\right] \quad j = 1, \dots, n-1. \tag{3.3}$$

The set of integer numbers  $n$  satisfying (3.3) for fixed  $p_0 \in \mathbb{R}$  may be described by the following construction (see, e.g., [TS], [KR1], [KR2]).

Let us define a sequence of real numbers  $p_i$  and sequences of integer numbers  $v_i, m_i, y_i$ :

$$p_0 = \frac{\pi}{\theta}, p_1 = 1, v_i = \left[\frac{p_i}{p_{i+1}}\right], p_{i+1} = p_{i-1} - v_{i-1}p_i \quad i = 1, 2, \dots \tag{3.4}$$

$$y_{-1} = 0, y_0 = 1, y_1 = v_0, y_{i+1} = y_{i-1} + v_i y_i \quad i = 0, 1, 2, \dots \tag{3.5}$$

$$m_0 = 0, m_1 = v_0, m_{i+1} = m_i + v_i \quad i = 0, 1, 2, \dots \tag{3.6}$$

It is clear that integer numbers  $v_i$  define the decomposition of  $p_0$  into a continuous fraction

$$p_0 = [v_0, v_1, v_2 \dots].$$

Let us define a piecewise linear function  $n_t, t \geq 1$

$$n_t = y_{i-1} + (t - m_i)y_i \quad \text{if } m_i \leq t < m_{i+1}.$$

Then for any integer  $n > 1$  there exists the unique rational number  $t$  such that  $n = n_t$ .

*Lemma 3.1.* [KR1]. The integer number  $n > 1$  satisfies (3.3) if and only if there exists an integer number  $t$  such that  $n = n_t$ .

We have two types of length 1 strings:

$$x_\alpha^1 \text{ with parity } v_1 = +1$$

$$x_\alpha^{m_1} \text{ with parity } v_{m_1} = -1.$$

All others strings have a length  $n = n_j$ , for some integer  $j$ , and parity

$$v_j = v_{n_j} = \exp\left(\pi i \left[\frac{n_j-1}{p_0}\right]\right).$$

Let us assume that all spins  $s_i$  have the following form:

$$2s_i = n_{\chi_i} - 1 \quad \chi_i \in \mathbb{Z}_+. \tag{3.7}$$

From the assumptions (3.1), (3.3) and (3.7) about spins, length and parity, a simple expression for the sums  $\theta_{n,m}(\infty)$  and  $\Phi_{n,2s}(\infty)$  follows.

In our paper we consider a special case of rational  $p_0$ . The case of irrational  $p_0$  may be obtained as a limit. So, we assume that  $p_0 = u/v \in \mathbb{Q}$ ,  $p_0 = [v_0, \dots, v_\alpha]$ ,  $v_0 \geq 2$ ,  $v_\alpha \geq 2$ . Furthermore, we assume that all strings have a length not greater than  $u$  (see [TS]). Therefore, for numbers  $p_i, v_i, y_i, m_i$  (see (3.4)–(3.6)), it is enough to keep only the indices  $i \leq \alpha + 1$ . We have also

$$p_{\alpha+1} = \frac{1}{y_\alpha} \quad p_0 = \frac{y_{\alpha+1}}{y_\alpha} \quad \text{and} \quad \text{GCD}(y_\alpha, y_{\alpha+1}) = 1.$$

Now we will state the results of calculations for the sums  $\theta_{n,m}(\infty)$  and  $\Phi_{n,m}(\infty)$ . Let us introduce

$$\begin{aligned} q_j &= (-1)^i (p_i - (j - m_i) p_{i+1}) && \text{if } m_i \leq j < m_{i+1} \\ r(j) &= i && \text{if } m_i \leq j < m_{i+1} \\ b_{jk} &= \frac{(-1)^{i-1}}{p_0} (q_k n_j - q_j n_k) && \text{if } n_j < n_k \\ b_{j,m_{i+1}} &= 1 && m_i < j < m_{i+1} \\ \theta_{j,k} &= \theta_{n_j, n_k}(\infty) && \Phi_{j,2s} = \Phi_{n_j, 2s}(\infty). \end{aligned} \tag{3.8}$$

*Theorem 3.2.* (Calculation of the sums  $\theta_{j,k}, \Phi_{j,2s}$ ) [KR1].

(1) If  $k > j$ , and  $(j, k) \neq (m_{\alpha+1} - 1, m_{\alpha+1})$  then

$$\theta_{jk} = -n_j \frac{q_k}{p_0}.$$

(2) If  $j = m_{\alpha+1} - 1, k = m_{\alpha+1}$ , then

$$\theta_{jk} = -n_j \frac{q_k}{p_0} + \frac{(-1)^{\alpha+1}}{2} = (-1)^\alpha \frac{p_0 - 2}{2p_0}.$$

(3) If  $1 \leq j \leq m_{\alpha+1}$ , then

$$\theta_{jj} = -n_j \frac{q_j}{p_0} + \frac{(-1)^{r(j)}}{2}.$$

(4) (Symmetry). For all  $1 \leq j, k \leq m_{\alpha+1}$

$$\theta_{jk} = \theta_{kj}.$$

(5) If  $2s = n_\chi - 1$ , then

$$\Phi_{k,2s} = \begin{cases} \frac{1}{2p_0} (q_k - q_k n_\chi) & \text{if } n_k > 2s \\ \frac{1}{2p_0} (q_k - q_\chi n_k) + \frac{(-1)^{r(k)-1}}{2} & \text{if } n_k \leq 2s. \end{cases}$$

Note that if  $p_0 = \nu_0$  is an integer then

$$2\Phi_{k,2s} = \begin{cases} \frac{2sk}{p_0} - \min(k, 2s) & \text{if } 1 \leq k, 2s + 1 < \nu_0 \\ 0 & \text{if } 1 \leq k \leq \nu_0, 2s + 1 > \nu_0. \end{cases}$$

Now we are going to calculate the vacancy numbers. By definition the vacancy numbers are equal to

$$P_{n_j}(\lambda) = 2Q_{\max}^{n_j} - \lambda_j + 1$$

where

$$Q_{\max}^{n_j} = (-1)^{i-1} \left( Q_\infty^{n_j} - \theta_{jj} - \frac{n_j}{2} \left\{ \frac{\sum 2s_m N_m - 2l}{p_0} \right\} \right) - \frac{1}{2} \quad m_i \leq j < m_{i+1}$$

and  $\{x\}$  is the fractional part of the real number  $x$ .

Here we put

$$Q_\infty^{n_j} = \sum_k \Phi_{n_j, 2s_k} N_k - \sum_k \theta_{n_j n_k} \lambda_k + \theta_{n_j n_j}.$$

Let us say a few words about our definition of the vacancy numbers  $P_{n_j}$ . In contrast with the  $XXZ$  model situation, it happens that the vector  $x = (\infty, \dots, \infty)$  for the  $XXZ$  case does not appear to be a formal solution to the Bethe equations (2.1). Another difficulty appears in finding a correct boundary for quantum numbers  $Q_\alpha^n$  (see (2.6)). A natural boundary is  $Q_\infty^{n_j}$  but this number does not appear to be an integer or half-integer one in general. Our choice is based on the attempt to have a combinatorial completeness of Bethe's states and some analytical considerations. In the following we will use the notation  $P_j(\lambda)$ ,  $Q_\infty^j$ ,  $Q_{\max}^j$ ,  $\dots$  instead of  $P_{n_j}(\lambda)$ ,  $Q_\infty^{n_j}$ ,  $Q_{\max}^{n_j}$ ,  $\dots$

After tedious calculations one can find

$$\begin{aligned} P_j(\lambda) &= a_j + 2 \sum_{k>j} b_{jk} \lambda_k & j \neq m_{\alpha+1} - 1, m_{\alpha+1} \\ P_{m_{\alpha+1}-1}(\lambda) &= a_{m_{\alpha+1}-1} + \lambda_{m_{\alpha+1}} \\ P_{m_{\alpha+1}}(\lambda) &= a_{m_{\alpha+1}} + \lambda_{m_{\alpha+1}-1} \end{aligned} \quad (3.9)$$

where

$$a_j = (-1)^{i-1} \left( \sum_m 2\Phi_{j,2s_m} \cdot N_m + \frac{2lq_j}{p_0} - n_j \left\{ \frac{\sum 2s_m N_m - 2l}{p_0} \right\} \right)$$

and  $b_{jk}$  for  $n_j < n_k$  are defined in (3.8).

From the string conjecture (see [TS], [KR2]) it follows that the number of Bethe's vectors with configuration  $\{\lambda_k\}$  is equal to

$$Z(N, s|\{\lambda_k\}) = \prod_j \binom{P_j(\lambda) + \lambda_j}{\lambda_j}.$$

The number of Bethe's vectors with fixed  $l$  is equal to

$$Z(N, s|l) = \sum_{\{\lambda_k\}} Z(N, s|\{\lambda_k\}) \quad (3.10)$$

where summation is taken over all configurations  $\{\lambda_k\}$ , such that  $\lambda_k \geq 0$ , and

$$\sum_{k=1}^{m_{\alpha+1}} n_k \lambda_k = l. \quad (3.11)$$

So, the total number of Bethe's vectors is equal to

$$Z = Z(N, s) = \sum_l Z(N, s|l) \quad (3.12)$$

where we assume that

$$Z(N, s|l) := Z(N, s|\sum 2s_m N_m - l) \quad \text{for } l \geq \sum s_m N_m.$$

The conjecture about combinatorial completeness of Bethe's states for the  $XXZ$  model means that

$$Z = \prod_m (2s_m + 1)^{N_m}. \quad (3.13)$$

#### 4. The main combinatorial identity

Let  $a_0 = 0, a_1, a_2, \dots, a_{m_{\alpha+1}}$  be a sequence of real numbers. Then we shall define inductively a sequence  $b_2, \dots, b_{m_1-1}, b_{m_1+1}, \dots, b_{m_{\alpha+1}}, b_{m_{\alpha+2}}$  by the following rules:

$$\begin{aligned} b_k &= 2a_{k-1} - a_{k-2} - a_k && \text{if } k \neq m_i, k \geq 2 \\ b_{m_i+1} &= 2a_{m_i-1} - a_{m_i-2} - a_{m_i+1} && \text{if } 1 \leq i \leq \alpha \\ b_{m_{\alpha+2}} &= a_{m_{\alpha+1}-1} - a_{m_{\alpha+1}-2} + a_{m_{\alpha+1}}. \end{aligned}$$

Then one can check that the converse formulae are

$$a_j = (-1)^{r(j)} \left( \frac{n_j}{p_0} q_{m_{\alpha+1}} (a_{m-1} - a_m) - 2 \sum_k \Phi_{jk} \cdot b_k \right)$$

where  $\Phi_{jk}$  were defined in (2.5).

For a given configuration  $\{\lambda_n\} = \lambda$  let us define the vacancy numbers

$$\begin{aligned} P_j(\lambda) &= a_j + 2 \sum_{k>j} b_{jk} \lambda_k && j \neq m_{\alpha+1} - 1, m_{\alpha+1} \\ P_{m_{\alpha+1}-1}(\lambda) &= a_{m_{\alpha+1}-1} + \lambda_{m_{\alpha+1}} \\ P_{m_{\alpha+1}}(\lambda) &= a_{m_{\alpha+1}} + \lambda_{m_{\alpha+1}-1}. \end{aligned}$$

Let us put

$$Z(\{a_k\} | l) = \sum_{\{\lambda_k\}} \prod_{k=1}^{m_{\alpha+1}} \binom{P_k(\lambda) + \lambda_k}{\lambda_k}$$

where summation is taken over all configurations  $\{\lambda_k\}$  such that

$$\sum_{k=1}^m n_k \lambda_k = l.$$

Note that a binomial coefficient  $\binom{\alpha}{\nu}$  for real  $\alpha$  and integer positive  $\nu$  is defined as

$$\binom{\alpha}{\nu} = \frac{\alpha(\alpha-1)\dots(\alpha-\nu+1)}{\nu!}.$$

*Theorem 4.1.* (The main combinatorial identity.) We have

$$Z(\{a_k\} | l) = \text{Res}_{u=0} f(u) u^{-l-1} du$$

where

$$\begin{aligned} f(u) &= (1+u)^{2l+2a_1-a_2} \prod_{k \neq m_i} \left( \frac{1-u^{n_k}}{1-u} \right)^{2a_{k-1}-a_k-a_{k-2}} \\ &\cdot \prod_{i=1}^{\alpha} \left( \frac{1-u^{y_i}}{1-u} \right)^{2a_{m_i-1}-a_{m_i-2}-a_{m_i+1}} \left( \frac{1-u^{y_{\alpha+1}}}{1-u} \right)^{a_{m_{\alpha+1}}+a_{m_{\alpha+1}-1}-a_{m_{\alpha+1}-2}}. \end{aligned}$$



*Proof.* We shall divide the proof into a few steps.

Step I. Let us put  $m_{\alpha+1} = m$ . We define a sequence of formal power series  $\varphi_1, \dots, \varphi_m$  in variables  $z_1, \dots, z_m, z_0$  by the following rules:

$$\begin{aligned} \varphi_m(z_m) &= (1 - z_m)^{-(a_m+1)}(1 - z_0(1 - z_m)^{-1})^{-1} \\ \varphi_{m-1}(z_{m-1}, z_m) &= (1 - z_{m-1})^{-(a_{m-1}+1)}\varphi_m((1 - z_{m-1})^{-1}z_m) \\ &\vdots \\ \varphi_k(z_k, \dots, z_m) &= (1 - z_k)^{-(a_k+1)}\varphi_{k+1}((1 - z_k)^{-2b_{k,k+1}} \\ &\quad \times z_{k+1}, \dots, (1 - z_k)^{-2b_{k,l}}z_l, \dots, (1 - z_k)^{-2b_{k,m}}z_m) \\ &\vdots \\ \varphi_1(z_1, \dots, z_m) &= (1 - z_1)^{-(a_1+1)}\varphi_2((1 - z_1)^{-2b_{1,2}}z_2, \dots, (1 - z_1)^{-2b_{1,l}} \\ &\quad \times z_l, \dots, (1 - z_1)^{-2b_{1,m}}z_m). \end{aligned}$$

*Lemma 4.2.* In the power series  $\varphi_1(z_1, \dots, z_m)$  a coefficient before  $z_0^{v_0}z_1^{v_1} \dots z_m^{v_m}$  is equal to

$$\prod_{j=1}^{m-1} \binom{P_j(v) + v_j}{v_j} \cdot \binom{a_m + v_m + v_0}{v_m}.$$

*Proof.*

$$\varphi_m(z_m) = \sum_{v_0, v_m} z_0^{v_0} z_m^{v_m} \binom{a_m + v_m + v_0}{v_m}.$$

Let us assume that

$$\varphi_k(z_k, \dots, z_m) = \sum_{v_0, v_k, \dots, v_m} A_k(v_k, \dots, v_m; v_0) z_0^{v_0} z_k^{v_k} \dots z_m^{v_m}$$

then

$$\begin{aligned} \varphi_{k-1}(z_{k-1}, \dots, z_m) &= (1 - z_{k-1})^{-(a_{k-1}+1)}\varphi_k((1 - z_k)^{-2b_{k,k+1}}z_{k+1}, \dots, (1 - z_k)^{-2b_{k,m}}z_m) \\ &= \sum_{v_0, v_k, \dots, v_m} A_k(v_k, \dots, v_m; v_0)(1 - z_{k-1})^{-(P_{k-1}(v)+1)}z_0^{v_0}z_k^{v_k} \dots z_m^{v_m} \\ &= \sum_{v_0, v_{k-1}, \dots, v_m} A_k(v_k, \dots, v_m; v_0) \binom{P_{k-1}(v) + v_{k-1}}{v_{k-1}} z_0^{v_0} z_{k-1}^{v_{k-1}} \dots z_m^{v_m}. \end{aligned}$$

Consequently,

$$A_{k-1}(v_{k-1}, v_k, \dots, v_m; v_0) = A_k(v_k, \dots, v_m; v_0) \cdot (P_{k-1}(v) + v_{k-1}v_{k-1}).$$

□

From lemma 4.2 it follows that the sum  $Z(\{a\}|l)$  is equal to the coefficient before  $t^l$  in the power series of  $\psi(z, t)$ , which has been obtained from  $\varphi_1(z_1, \dots, z_m)$  after substitution

$$\begin{aligned} z_j &= t^{n_j} \quad j \neq m - 1 \\ z_{m-1} &= t^{n_{m-1}}z_0^{-1}. \end{aligned}$$

Step II. Calculation of the power series for  $\psi(z, t)$ . Let us define

$$\begin{aligned} z_k^{(l)} &:= (1 - z_l^{(l-1)})^{-2b_{l,k}} \cdot z_k^{(l-1)} \quad l \geq 1 \\ z_k^{(0)} &= t^{n_k} \quad \text{if } k \neq m - 1 \text{ and } z_{m-1}^{(0)} = t^{n_{m-1}}z_0^{-1}. \end{aligned} \tag{4.1}$$

Then we have

$$\begin{aligned}
 \varphi_1(z_1, \dots, z_m) &= (1 - z_1)^{-(a_1+1)} \varphi_2(z_2^{(1)}, z_3^{(1)}, \dots, z_m^{(1)}) \\
 &= (1 - z_1)^{-(a_1+1)} (1 - z_2^{(1)})^{-(a_2+1)} \varphi_3(z_3^{(2)}, z_4^{(2)}, \dots, z_m^{(2)}) \\
 &\quad \vdots \\
 &= \prod_{j=1}^{m-1} (1 - z_j^{(j-1)})^{-(a_j+1)} \cdot \varphi_{m-1}(z_{m-1}^{(m-2)}, z_m^{(m-2)}).
 \end{aligned} \tag{4.2}$$

In order to compute a formal series  $z_k^{(l)}$ , we define (see, e.g., [K1]) a sequence of polynomials  $Q_m(t)$  using the following recurrence relation:

$$\begin{aligned}
 Q_{m+1}(t) &= Q_m(t) - t Q_{m-1}(t) & m \geq 0 \\
 Q_0(t) &= Q_{-1}(t) = 1.
 \end{aligned}$$

*Lemma 4.3.* (Formulae for power series  $z_k^{(l)}$ .) Let us assume that  $m_i \leq k < m_{i+1}$  and put  $m_0 := 1$ . Then we have ( $Q_k := Q_k(t)$ )

- (1)  $z_k^{(k-1)} = Q_{k-1}^{-2} Q_{m_i-2} z_k^{(0)}$ .
- (2)  $1 - z_k^{(k-1)} = Q_k Q_{k-1}^{-2} Q_{k-2}$ , if  $k \neq m_i$ .
- (3) If  $k = m_i$ ,  $i \geq 1$ , then  $1 - z_k^{(k-1)} = Q_k Q_{k-1}^{-2} Q_{m_{i-1}-2}$ .
- (4) After specialization  $t := u/(1+u)^2$  one can find (note that  $m_i \leq k < m_{i+1}$ )

$$Q_k(u) = 1 - \frac{1 - u^{n_k+2y_i}}{(1-u)(1+u)^{n_k+2y_i-1}}.$$

- (5) If  $k \neq m_i + 1$  and  $m_i \leq k < m_{i+1}$ , then

$$z_k^{(k-2)} = Q_{k-3}^2 Q_{k-2}^{-4} Q_{m_i-2}^2 z_k^{(0)}.$$

*Proof.* This follows by induction from (4.1) and the properties of polynomials  $Q_k(t)$  (compare [K1], lemma 2).  $\square$

*Corollary 4.4.* (1)

$$z_m^{(m-2)} = Q_{m-3}^2 Q_{m-2}^{-2} t^{n_m} \quad z_{m-1}^{(m-2)} = Q_{m-2}^{-2} Q_{m_{\alpha-2}}^2 t^{n_{m-1}} z_0^{-1}.$$

(2) Let us denote by  $\varphi_{m-1}(u, z_0)$  a specialization  $t = u/(1-u)^2$  of formal series  $\varphi_{m-1}(z_{m-1}^{(m-2)}, z_m^{(m-1)})$  and let  $\varphi_{m-1}(u)$  be a constant term of series  $\varphi_{m-1}(u, z_0)$  with respect to variable  $z_0$ . Then

$$\varphi_{m-1}(u) = (1 - u^{y_{\alpha+1}})^{a_m + a_{m-1} + 1} (1 - u^{y_{\alpha+1} - y_{\alpha}})^{-(a_{m-1} + 1)} (1 - u^{y_{\alpha}})^{-(a_m + 1)}.$$

Note that  $m = m_{\alpha+1}$ .

Step III. Combining (4.2), lemma 4.3 and corollary 4.4 after some simplifications we obtain a proof of theorem 4.1.  $\square$

*Corollary 4.5.* (Combinatorial completeness of Bethe's states for XXZ model of arbitrary spins.)

$$Z = \prod_m (2s_m + 1)^{N_m}. \tag{4.3}$$

Examples below give an illustration to our result about completeness of Bethe's states for the spin- $\frac{1}{2}$  XXZ model (examples 1 and 3) and for the spin-1 XXZ model (example 4).

*Example 1.* We compute firstly the quantities  $q_j, a_j$  (see (3.8)) and after this consider a numerical example. From (3.4)–(3.6) and (3.8) it follows that

$$q_j = (-1)^i \frac{p_0 - n_j p_{i+1}}{y_i}.$$

Using theorem 3.2(5) we obtain (see (3.9))

$$\begin{aligned} a_j = & (-1)^{i-1} n_j \left[ \frac{\sum 2s_m N_m - 2l}{p_0} \right] + (-1)^i (n_j + q_j) \left( \frac{\sum 2s_m N_m - 2l}{p_0} \right) \\ & + \frac{n_j}{p_0} \sum_{\{m: 2s_m \leq n_j\}} N_m \left( \frac{p_{i+1}}{y_i} (2s_m + 1) + (-1)^i q_\chi \right) \\ & + \sum_{\{m: 2s_m \leq n_j\}} N_m \left( 1 - \frac{1}{y_i} (2s_m + 1) \right). \end{aligned} \quad (4.4)$$

Let us consider the case when all spins are equal to  $\frac{1}{2}$  and let  $N$  be the number of spins, then

(i)  $0 \leq j < m_1 (= v_0)$ . Then  $r(j) = i = 0$  and  $n_j = j, q_j = p_0 - j$ ,

$$a_j = -n_j \left[ \frac{N - 2l}{p_0} \right] + N - 2l + \delta_{n_j, 1} \frac{N}{p_0} (2 - p_0 + q_\chi).$$

(ii)  $m_1 \leq j < m_2 (= v_0 + v_1)$ . Then  $r(j) = 1$  and  $n_j = 1 + (j - m_1)v_0$ ,  $q_j = (p_0 - v_0)(j - m_1) - 1$ ,

$$a_j = n_j \left[ \frac{N - 2l}{p_0} \right] - \frac{N - 2l}{v_0} (n_j - 1) - \delta_{n_j, 1} \frac{N}{p_0} (2 - p_0 + q_\chi).$$

For example,

$$a_{m_1} = \left[ \frac{N - 2l}{p_0} \right] - \frac{N}{p_0} (2 - p_0 + q_\chi).$$

(iii)  $m_2 \leq j < m_3 (= v_0 + v_1 + v_2)$ . Then  $r(j) = 2$  and  $n_j = v_0 + (j - m_2)(1 + v_0 v_1)$ ,  $q_j = p_0 - v_0 - (j - m_2)(1 - v_1(p_0 - v_0))$

$$a_j = -n_j \left[ \frac{N - 2l}{p_0} \right] + \frac{N - 2l}{v_0 + (1/v_1)} \left( n_j + \frac{1}{v_1} \right).$$

Consequently,

$$a_{m_2} = -v_0 \left[ \frac{N - 2l}{p_0} \right] + (N - 2l).$$

Now let us assume  $p_0 = 3 + \frac{1}{3}, N = 5$ . It is clear that in our case  $\chi = 2$  (see (3.7)) and  $q_\chi = p_0 - 2$ . Below we give all solutions  $\lambda = \{\lambda_1, \lambda_2, \dots\}$  to the equation (3.11) when  $0 \leq l \leq 2$  and compute the corresponding vacancy numbers  $P_j = P_j(\lambda)$  (see (3.9)) and

number of states  $Z = Z(N, \frac{1}{2}|\{\lambda_k\})$  (see (3.10) and (3.12)):

$$\begin{array}{llll}
 l = 0 & \{0\} & P_j = 0 & Z = 1 \\
 & & & Z(5, \frac{1}{2}|0) = 1 \\
 l = 1 & \{1, 0, 0\} & P_1 = 3 & Z = 4 \\
 & \{0, 0, 1\} & P_3 = 0 & Z = 1 \\
 & & & Z(5, \frac{1}{2}|1) = 5 \\
 l = 2 & \{0, 1, 0\} & P_2 = 1 & Z = 2 \\
 & \{2, 0, 0\} & P_1 = 1 & Z = 3 \\
 & \{0, 0, 2\} & P_3 = 0 & Z = 1 \\
 & \{1, 0, 1\} & \begin{cases} P_1 = 3 \\ P_3 = 0 \end{cases} & Z = 4 \\
 & & & Z(5, \frac{1}{2}|2) = 10.
 \end{array}$$

Consequently,

$$Z(N = 5, \frac{1}{2}) = 2(1 + 5 + 10) = 32 = 2^5.$$

Note that our formula (3.10) for the number of Bethe's states with fixed spin  $l$ , namely  $Z(N, s|l)$ , works for  $l \geq \sum s_m N_m$  as well as for small  $l \leq \sum s_m N_m$ .

In the appendix we consider two additional examples, one when all spins are equal to  $\frac{1}{2}$ , another when all spins are equal to 1. The last example seems to be interesting because a non-admissible configuration appears.

*Remark 1.* It is easy to see that for fixed  $l$  and sufficiently big  $N = \sum 2s_m N_m$  all vacancy numbers  $P_j(\lambda)$  are non-negative. This is not the case for particular  $N$  and we must consider really the configurations with

$$P_j(\lambda) + \lambda_j < 0 \quad \text{for some } j \tag{4.5}$$

in order to have a correct answer for  $Z^{XXZ}(N, s|l)$ . See the appendix, example 4,  $l = 4$ , ( $\clubsuit$ ). Let us note that for the XXX model the non-admissible configurations (i.e. those satisfying (4.5)) give a zero contribution to the sum  $Z^{XXX}(N, s|l)$  [K2].

*Remark 2.* One can rewrite the expressions (3.9) for vacancy numbers in the following form if  $m_i \leq j < m_{i+1}$ ,

$$\begin{aligned}
 P_j(\lambda) = & (-1)^{i-1} \left( \sum_m 2\Phi_{j,2s_m} \cdot N_m - n_j \left\{ \frac{\sum 2s_m N_m - 2l}{p_0} \right\} \right) \\
 & - \sum_k 2(-1)^{r(k)} \tilde{\theta}_{jk} \lambda_k - \delta_{j,m_{\alpha+1}-1} \lambda_{m_{\alpha+1}} + \delta_{j,m_{\alpha+1}} \lambda_{m_{\alpha+1}-1}
 \end{aligned}$$

where  $\tilde{\theta}_{jk} = (-1)^{r(j)+r(k)} n_j q_k / p_0$ , if  $j \leq k$  and  $\tilde{\theta}_{jk} = \tilde{\theta}_{kj}$ .

Let us introduce the symmetric matrix  $\Theta = (\tilde{\theta}_{ij})_{1 \leq i, j \leq m_{\alpha+1}}$ . We can find the inverse matrix  $\Theta^{-1} := (c_{ij})$  and compute its determinant.

*Theorem 4.6.* Matrix elements  $c_{ij}$  of the inverse matrix  $\Theta^{-1}$  are given by the following rules:

- (i)  $c_{ij} = c_{ji}$  and  $c_{ij} = 0$ , if  $|i - j| \geq 2$ ;
- (ii)  $c_{j-1,j} = (-1)^{i-1}$ , if  $m_i \leq j < m_{i+1}$ ;



The thermodynamical limit of (5.1) (i.e.  $N_m \rightarrow \infty$ ) comes to

$$\sum_{\lambda} \frac{q^{\frac{1}{2}\tilde{\lambda}B\tilde{\lambda}'}}{\prod_j (q)_{\lambda_j}}. \tag{5.2}$$

Summation in (5.2) is the same as in (5.1) and  $(q)_n := (1-q) \cdots (1-q^n)$ . Here  $B = C_1 \otimes \Theta$  and  $C_1 = (2)$  is the Cartan matrix of type  $A_1$ .

It is an interesting problem to find a representation theory meaning of (5.2), when  $B = C_k \otimes \Theta$  and  $C_k$  is the Cartan matrix of type  $A_k$ .

Another interesting question concerns the degeneration of Bethe's states for the XXZ model into those for the XXX one. More exactly, we had proved (see (4.3)) that

$$\prod_m (2s_m + 1)^{N_m} = \sum_{l=0}^N Z^{XXZ}(N, s|l) \tag{5.3}$$

where  $N = \sum_m 2s_m N_m$  and  $Z^{XXZ}(N, s|l)$  is given by (3.10).

On the other hand, it follows from a combinatorial completeness of Bethe's states for the XXX model (see [K1]) that

$$\prod_m (2s_m + 1)^{N_m} = \sum_{l \geq 0}^{\frac{1}{2}N} (N - 2l + 1) Z^{XXX}(N, s|l) \tag{5.4}$$

where  $Z^{XXX}(N, s|l)$  is the multiplicity of the  $(\frac{1}{2}N - l)$ -spin irreducible representation  $V_{\frac{1}{2}N-l}$  of  $sl(2)$  in the tensor product

$$V_{s_1}^{\otimes N_1} \otimes \cdots \otimes V_{s_m}^{\otimes N_m}.$$

It is an interesting question to find a combinatorial proof that

$$\text{RHS}(5.3) = \text{RHS}(5.4).$$

Another interesting task is to compare our results with those obtained in [KM]. We intend to consider these questions and also to study in more detail the case  $p_0 = \nu_0$  as an integer and all spins equal to  $(\nu_0 - 2)/2$  in separate publications.

**Acknowledgments**

We are pleased to thank for hospitality our colleagues from Tokyo University, where this work was completed.

**Appendix**

*Example 3.* Using the same notation as in example 1, we consider the case  $s = \frac{1}{2}$ ,  $p_0 = 3 + \frac{1}{3}$ ,  $N = 8$  and compute the vacancy numbers  $P_j(\lambda)$  and numbers of states  $Z = Z(N, \frac{1}{2}|\{\lambda_k\})$ :

$l = 0$	$\{0\}$	$P_j = 0$	$Z = 1$	$Z(8, \frac{1}{2} 0) = 1$
$l = 1$	$\{1, 0, 0\}$ $\{0, 0, 1\}$	$P_1 = 5$ $P_3 = 1$	$Z = 6$ $Z = 2$	$Z(8, \frac{1}{2} 1) = 8$
$l = 2$	$\{2, 0, 0\}$ $\{0, 0, 2\}$ $\{0, 1, 0\}$ $\{1, 0, 1\}$	$P_1 = 3$ $P_3 = 1$ $P_2 = 2$ $\left\{ \begin{array}{l} P_1 = 5 \\ P_3 = 1 \end{array} \right.$	$Z = 10$ $Z = 3$ $Z = 3$ $Z = 12$	$Z(8, \frac{1}{2} 2) = 28$
$l = 3$	$\{3, 0, 0\}$ $\{0, 0, 3\}$ $\{0, 0, 0, 0, 0, 1\}$ $\{1, 1, 0\}$ $\{0, 1, 1\}$ $\{2, 0, 1\}$ $\{1, 0, 2\}$	$P_1 = 2$ $P_3 = 0$ $P_6 = 2$ $\left\{ \begin{array}{l} P_1 = 4 \\ P_2 = 2 \end{array} \right.$ $\left\{ \begin{array}{l} P_2 = 4 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 4 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 6 \\ P_3 = 0 \end{array} \right.$	$Z = 10$ $Z = 1$ $Z = 3$ $Z = 15$ $Z = 5$ $Z = 15$ $Z = 7$	$Z(8, \frac{1}{2} 3) = 56$
$l = 4$	$\{4, 0, 0\}$ $\{0, 0, 4\}$ $\{0, 2, 0\}$ $\{0, 0, 0, 1\}$ $\{2, 1, 0\}$ $\{0, 1, 2\}$ $\{3, 0, 1\}$ $\{1, 0, 3\}$ $\{2, 0, 2\}$ $\{1, 0, 0, 0, 0, 1\}$ $\{0, 0, 1, 0, 0, 1\}$	$P_1 = 0$ $P_3 = 0$ $P_2 = 0$ $P_4 = 0$ $\left\{ \begin{array}{l} P_1 = 2 \\ P_2 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_2 = 4 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 2 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 6 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 4 \\ P_3 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_1 = 4 \\ P_6 = 0 \end{array} \right.$ $\left\{ \begin{array}{l} P_3 = 2 \\ P_6 = 0 \end{array} \right.$	$Z = 1$ $Z = 1$ $Z = 1$ $Z = 1$ $Z = 6$ $Z = 5$ $Z = 10$ $Z = 7$ $Z = 15$ $Z = 5$ $Z = 3$	

$$\{1, 1, 1, 0, 0, 0\} \begin{cases} P_1 = 4 \\ P_2 = 2 \\ P_3 = 0 \end{cases} \quad Z = 15$$

$$Z(8, \frac{1}{2}|4) = 70.$$

Consequently,

$$Z(N = 8, \frac{1}{2}|l) = \binom{8}{l} \quad 0 \leq l \leq 4$$

and

$$Z(N = 8, \frac{1}{2}) = 2(1 + 8 + 28 + 56) + 70 = 256 = 2^8.$$

*Example 4.* Let us consider the case when all spins are equal to 1 and let  $N$  be the number of spins. We compute firstly the quantities  $a_j$  (see (3.8)) and after this consider a numerical example.

(i)  $0 \leq j < m_1 (= v_0)$ . Then  $r(j) = i = 0$  and  $n_j = j, q_j = p_0 - j$ ,

$$a_j = \begin{cases} -j \left[ \frac{2N - 2l}{p_0} \right] + 2N - 2l & \text{if } j > 2 \\ -j \left[ \frac{2N - 2l}{p_0} \right] + \frac{jN}{p_0} (3 + q_\chi) - 2l & \text{if } j \leq 2. \end{cases}$$

(ii)  $m_1 \leq j < m_2 (= v_0 + v_1)$ . Then  $r(j) = 1$  and  $n_j = 1 + (j - m_1)v_0, q_j = (p_0 - v_0)(j - m_1) - 1$ ,

$$a_j = n_j \left[ \frac{2N - 2l}{p_0} \right] - \frac{2N - 2l}{v_0} (n_j - 1) - \delta_{n_j,1} \frac{N}{p_0} (3 - p_0 + q_\chi).$$

(iii)  $m_2 \leq j < m_3 (= v_0 + v_1 + v_2)$ . Then  $r(j) = 2$  and  $n_j = v_0 + (j - m_2)(1 + v_0v_1), q_j = p_0 - v_0 - (j - m_2)(1 - v_1(p_0 - v_0))$ ,

$$a_{m_2} = -v_0 \left[ \frac{2N - 2l}{p_0} \right] + 2N - 2l.$$

Now let us assume  $p_0 = 3 + \frac{1}{3}, N = 5$ . It is clear that  $\chi = 6$  and  $q_\chi = \frac{1}{3}$ . Below we give all solutions  $\lambda = \{\lambda_1, \lambda_2, \dots\}$  to the equation (3.11) when  $0 \leq l \leq 5$  and compute the corresponding vacancy numbers  $P_j = P_j(\lambda)$  (see (3.9)) and number of states  $Z = Z(N, 1|\{\lambda_k\})$  (see (3.10) and (3.12)):

$l = 0$	$\{0\}$	$P_j = 0$	$Z = 1$	
				$Z(5, 1 0) = 1$
$l = 1$	$\{1, 0, 0\}$	$P_1 = 1$	$Z = 2$	
	$\{0, 0, 1\}$	$P_3 = 2$	$Z = 3$	
				$Z(5, 1 1) = 5$
$l = 2$	$\{2, 0, 0\}$	$P_1 = 0$	$Z = 1$	
	$\{0, 0, 2\}$	$P_3 = 1$	$Z = 3$	
	$\{0, 1, 0\}$	$P_2 = 4$	$Z = 5$	
	$\{1, 0, 1\}$	$\begin{cases} P_1 = 2 \\ P_3 = 1 \end{cases}$	$Z = 6$	
				$Z(5, 1 2) = 15$



$l = 3$	$\{3, 0, 0\}$	$P_1 = -2$	$Z = 0$	$Z(5, 1 3) = 30$
	$\{0, 0, 3\}$	$P_3 = 1$	$Z = 4$	
	$\{0, 0, 0, 0, 0, 1\}$	$P_6 = 1$	$Z = 2$	
	$\{1, 1, 0\}$	$\begin{cases} P_1 = 0 \\ P_2 = 2 \end{cases}$	$Z = 3$	
	$\{0, 1, 1\}$	$\begin{cases} P_2 = 4 \\ P_3 = 1 \end{cases}$	$Z = 10$	
	$\{2, 0, 1\}$	$\begin{cases} P_1 = 0 \\ P_3 = 1 \end{cases}$	$Z = 2$	
	$\{1, 0, 2\}$	$\begin{cases} P_1 = 2 \\ P_3 = 1 \end{cases}$	$Z = 9$	
$l = 4$	$\{4, 0, 0\}$	$P_1 = -3$	$Z = 0$	$(\clubsuit)$
	$\{0, 0, 4\}$	$P_3 = 0$	$Z = 1$	
	$\{0, 2, 0\}$	$P_2 = 2$	$Z = 6$	
	$\{0, 0, 0, 1\}$	$P_4 = -2$	$Z = -1$	
	$\{2, 1, 0\}$	$\begin{cases} P_1 = -1 \\ P_2 = 2 \end{cases}$	$Z = 0$	
	$\{0, 1, 2\}$	$\begin{cases} P_2 = 6 \\ P_3 = 0 \end{cases}$	$Z = 7$	
	$\{3, 0, 1\}$	$\begin{cases} P_1 = -1 \\ P_3 = 0 \end{cases}$	$Z = 0$	
	$\{1, 0, 3\}$	$\begin{cases} P_1 = 3 \\ P_3 = 0 \end{cases}$	$Z = 4$	
	$\{2, 0, 2\}$	$\begin{cases} P_1 = 1 \\ P_3 = 0 \end{cases}$	$Z = 3$	
	$\{1, 0, 0, 0, 0, 1\}$	$\begin{cases} P_1 = 1 \\ P_6 = 2 \end{cases}$	$Z = 6$	
$l = 5$	$\{0, 0, 1, 0, 0, 1\}$	$\begin{cases} P_3 = 2 \\ P_6 = 2 \end{cases}$	$Z = 9$	$Z(5, 1 4) = 45$
	$\{1, 1, 1, 0, 0, 0\}$	$\begin{cases} P_1 = 1 \\ P_2 = 4 \\ P_3 = 0 \end{cases}$	$Z = 10$	
	$\{5, 0, 0\}$	$P_1 = -5$	$Z = 0$	
$l = 5$	$\{0, 0, 5\}$	$P_3 = 0$	$Z = 1$	
	$\{4, 0, 1\}$	$\begin{cases} P_1 = -3 \\ P_3 = 0 \end{cases}$	$Z = 0$	
	$\{1, 0, 4\}$	$\begin{cases} P_1 = 3 \\ P_3 = 0 \end{cases}$	$Z = 4$	

$$\begin{array}{l}
 \{3, 0, 2\} \quad \left\{ \begin{array}{l} P_1 = -1 \\ P_3 = 0 \end{array} \right. \quad Z = 0 \\
 \{2, 0, 3\} \quad \left\{ \begin{array}{l} P_1 = 1 \\ P_3 = 0 \end{array} \right. \quad Z = 3 \\
 \{3, 1, 0\} \quad \left\{ \begin{array}{l} P_1 = -3 \\ P_2 = 0 \end{array} \right. \quad Z = 0 \\
 \{0, 1, 3\} \quad \left\{ \begin{array}{l} P_2 = 6 \\ P_3 = 0 \end{array} \right. \quad Z = 7 \\
 \{1, 2, 0\} \quad \left\{ \begin{array}{l} P_1 = -1 \\ P_2 = 0 \end{array} \right. \quad Z = 0 \\
 \{0, 2, 1\} \quad \left\{ \begin{array}{l} P_2 = 2 \\ P_3 = 0 \end{array} \right. \quad Z = 6 \\
 \{1, 0, 0, 1\} \quad \left\{ \begin{array}{l} P_1 = 1 \\ P_4 = 0 \end{array} \right. \quad Z = 2 \\
 \{0, 0, 1, 1\} \quad \left\{ \begin{array}{l} P_3 = 2 \\ P_4 = 0 \end{array} \right. \quad Z = 3 \\
 \{0, 1, 0, 0, 0, 1\} \quad \left\{ \begin{array}{l} P_2 = 2 \\ P_6 = 0 \end{array} \right. \quad Z = 3 \\
 \{2, 0, 0, 0, 0, 1\} \quad \left\{ \begin{array}{l} P_1 = -1 \\ P_6 = 0 \end{array} \right. \quad Z = 0 \\
 \{0, 0, 2, 0, 0, 1\} \quad \left\{ \begin{array}{l} P_3 = 2 \\ P_6 = 0 \end{array} \right. \quad Z = 6 \\
 \{1, 0, 1, 0, 0, 1\} \quad \left\{ \begin{array}{l} P_1 = 1 \\ P_3 = 2 \\ P_6 = 0 \end{array} \right. \quad Z = 6 \\
 \{2, 1, 1\} \quad \left\{ \begin{array}{l} P_1 = -1 \\ P_2 = 2 \\ P_3 = 0 \end{array} \right. \quad Z = 0 \\
 \{1, 1, 2\} \quad \left\{ \begin{array}{l} P_1 = 1 \\ P_2 = 4 \\ P_3 = 0 \end{array} \right. \quad Z = 10
 \end{array}$$

$$Z(5, 1|5) = 51$$

$$Z(N = 5, 1) = 2(1 + 5 + 15 + 30 + 45) + 51 = 243 = 3^5.$$

### References

- [TS] Takahashi M and Suzuki M 1972 One-dimensional anisotropic Heisenberg model at finite temperatures *Prog. Theor. Phys.* **48** 2187–209
- [K1] Kirillov A N 1984 Combinatorial identities and completeness of states for the generalized Heisenberg magnet *Zap. Nauch. Sem. LOMI* **131** 88–105
- [K2] Kirillov A N 1988 On the Kostka–Green–Foulkes polynomials and Clebsch–Gordon numbers *J. Geom. Phys.* **5** 365–89

- [KR1] Kirillov A N and Reshetikhin N Yu 1985 Properties of kernels of integrable equations for  $XXZ$  model of arbitrary spin *Zap. Nauch. Sem. LOMI* **146** 47–91
- [KR2] Kirillov A N and Reshetikhin N Yu 1987 Exact solution of the integrable  $XXZ$  Heisenberg model with arbitrary spin *J. Phys. A: Math. Gen.* **20** 1565–97
- [FT] Faddeev L D and Takhtadjan L A 1981 Spectrum and scattering of excitations in one dimensional isotropic Heisenberg model *Zap. Nauch. Sem. LOMI* **109** 134
- [EKK] Essler F, Korepin V E and Schoutens K 1992 Fine structure of the Bethe ansatz for the spin- $\frac{1}{2}$  Heisenberg  $XXX$  model *J. Phys. A: Math. Gen.* **25** 4115–26
- [KM] Kedem R and McCoy B 1993 Construction of modular branching functions from Bethe's equations in the 3-state Potts chain *J. Stat. Phys.* **74** 865
- [BM] Bercovich A and McCoy B 1996 Continued fractions and fermionic representations for characters of  $\mu(p, p')$  minimal models *Lett. Math. Phys.* **37** 49–66